

STUDY OF CERTAIN NEW TENSORS IN A FINSLER SPACE OF THREE-DIMENSIONS

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Received: 26 Oct 2019

Accepted: 02 Nov 2019

Published: 30 Nov 2019

ABSTRACT

In one of his earlier papers in (1990), the author [7] has defined and studied several new tensors of second and third order. The author, further in Ref. [9] has defined and studied several properties of asymmetric third order new tensor D_{ijk} , which is similar to the Cartan's torsion tensor [1], C_{ijk} of a Finsler space of three dimensions. This tensor, however, satisfies $D_{ijk} l^i = 0$, $D_{ijk} g^{jk} = D_i = D n_i$. Based on this tensor, author has defined several other tensors including $Q_{ijk} = D_{ijk}/0$, similar to P_{ijk} and D_{ijkh} similar to third curvature tensor S_{ijkh} . The purpose of the present paper is to define and study a tensor U_{ijkh} , similar to second curvature tensor P_{ijkh} and V_{ijkh} similar to a very important tensor T_{ijkh} , which was introduced independently by Kawaguchi [3] and Matsumoto [5] in (1972).

KEYWORDS: Finsler Space F^3 , Tensors D_{ijk} , Q_{ijk} , U_{ijkh} and V_{ijkh}

INTRODUCTION

Let F^3 be a three-dimensional Finsler space with the Moor's frame (l_i, m_i, n_i) . Corresponding to this frame, the metric tensor, angular metric tensor and (h) hv-torsion tensors are given by Matsumoto [6] and Rund [10] as

$$g_{ij} = l_i l_j + m_i m_j + n_i n_j, \quad h_{ij} = m_i m_j + n_i n_j, \quad (1.1)$$

and

$$C_{ijk} = C_{(1)} m_i m_j m_k + C_{(2)} n_i n_j n_k + \sum_{(ijk)} \{ C_{(3)} m_i n_j n_k - C_{(2)} m_i m_j n_k \} \quad (1.2)$$

In a three-dimensional Finsler space F^3 , h- and v-covariant derivatives of the unit vectors l^i, m^i and n^i are defined as [6]

$$l^i_{/j} = 0, \quad m^i_{/j} = n^i h_j, \quad n^i_{/j} = -m^i h_j \quad (1.3)$$

and

$$l^i_{/j} = L^{-1} h^i_j, \quad m^i_{/j} = L^{-1} (-l^i m_j + n^i v_j), \quad n^i_{/j} = -L^{-1} (l^i n_j + m^i v_j) \quad (1.4)$$

where $v_j = v_{2)3\gamma} e_{\gamma j}$ and $h_j = H_{2)3\gamma} e_{\gamma j}$.

In, general for a tensor field K^h_m , we have

$$K^h_{m/r} = \partial_r K^h_m - N^j_r \Delta_j K^h_m + K^k_m F_k^h_r - K^h_k F^k_m_r \quad (1.5)$$

and

$$K^h_{m/r} = \Delta_r K^h_m + K^j_m C^h_{j r} - K^h_j C^j_{m r} \quad (1.6)$$

where $\partial_r = \partial/\partial x^r$ and $\Delta_r = \partial/\partial y^r$.

The second and third curvature tensors in the sense of E. Cartan [1] are given by

$$P_{ijkh} = \zeta_{(i,j)} \{A_{jkh/i} + A_{ikr} P^r_{jh}\} \quad (1.7)$$

$$S_{ijkh} = \zeta_{(h,k)} \{A_{ihr} A^r_{jk}\} \quad (1.8)$$

such that

$$\zeta_{(h,k)} \{P_{ijkh}\} = -S_{ijkh/0} \quad (1.9)$$

where $\zeta_{(h,k)} \{ \}$ means interchange of indices h and k and subtraction.

The tensor D_{ijk} was introduced and defined by Rastogi [9], such that it satisfies $D_{ijk} l^i = 0$, $D_{ijk} g^{jk} = D_i = D n_i$ and is given as

$$D_{ijk} = D_{(1)} m_i m_j m_k + D_{(2)} n_i n_j n_k + \sum_{(i,j,k)} \{D_{(3)} m_i m_j n_k + D_{(4)} m_i n_j n_k\} \quad (1.10)$$

where $D_{(1)}$, $D_{(2)}$, $D_{(3)}$ and $D_{(4)}$ are scalars satisfying,

$$D_{(2)} + D_{(3)} = D, D_{(1)} + D_{(4)} = 0 \quad (1.11)$$

Tensor U_{ijkh}

Corresponding to second curvature tensor of Cartan [1], P_{ijkh} , we here define the curvature tensor U_{ijkh} as follows:

$$U_{ijkh} = \zeta_{(i,j)} \{D_{jkh/i} + D_{ikr} Q^r_{jh}\} \quad (2.1)$$

From equation (1.10), we can get

$$D_{jkh/i} = m_j m_k m_h (D_{(1)/i} - 3 D_{(3)} h_i) + n_j n_k n_h (D_{(2)/i} - 3 D_{(1)} h_i) + \sum_{(j,k,h)} [m_j m_k n_h \{D_{(3)/i} + (D_{(1)} - 2 D_{(4)}) h_i\} - m_j n_k n_h \{D_{(1)/i} + (D_{(2)} - 2 D_{(3)}) h_i\}] \quad (2.2)$$

whereas $Q_{ijk} = D_{ijk/0}$, can be given as

$$Q_{ijk} = \{D_{(1)/0} - 3 D_{(3)} h_0\} m_i m_j m_k + \{(D_{(2)/0} - 3 D_{(1)} h_0) n_i n_j n_k\} + \sum_{(i,j,k)} [\{D_{(3)/0} + 3 D_{(1)} h_0\} m_i m_j n_k - \{D_{(1)/0} + (D_{(2)} - 2 D_{(3)}) h_0\} m_i n_j n_k] \quad (2.3)$$

From equation (2.1), by virtue of (2.2) and (2.3), we can obtain

$$U_{ijkh} = {}^1 A_{ij} m_k m_h + {}^2 A_{ij} m_k n_h + {}^3 A_{ij} m_h n_k + {}^4 A_{ij} n_h n_k, \quad (2.4)$$

where, we have assumed

$${}^1 A_i = D_{(1)/i} - 3 D_{(3)} h_i, {}^2 A_i = D_{(1)/i} + (D_{(2)} - 2 D_{(3)}) h_i, \quad (2.5) a$$

$${}^3 A_i = D_{(2)/i} - 3 D_{(1)} h_i, {}^4 A_i = D_{(3)/i} + 3 D_{(1)} h_i \quad (2.5) b$$

and

$${}^1 A_0 = {}^1 A_i l^i, {}^2 A_0 = {}^2 A_i l^i, {}^3 A_0 = {}^3 A_i l^i, {}^4 A_0 = {}^4 A_i l^i, \quad (2.5) c$$

such that

$${}^1A_{ij} = \zeta_{(i,j)} [{}^1A_i m_j + {}^4A_i n_j + \{D_{(3)} ({}^1A_0 - {}^2A_0) - (D_{(1)} - D_{(4)}) {}^4A_0\} m_j n_i] \tag{2.6a}$$

$${}^2A_{ij} = \zeta_{(i,j)} [{}^4A_i m_j - {}^2A_i n_j + \{D_{(3)} ({}^4A_0 - {}^3A_0) + (D_{(1)} - D_{(4)}) {}^2A_0\} m_j n_i] \tag{2.6b}$$

$${}^3A_{ij} = \zeta_{(i,j)} [{}^4A_i m_j - {}^2A_i n_j + \{D_{(4)} ({}^1A_0 - {}^2A_0) + (D_{(2)} - D_{(3)}) {}^4A_0\} m_j n_i] \tag{2.6c}$$

$${}^4A_{ij} = \zeta_{(i,j)} [{}^2A_j m_i - {}^3A_j n_i + \{D_{(4)} ({}^4A_0 - {}^3A_0) - (D_{(2)} - D_{(3)}) {}^2A_0\} m_j n_i] \tag{2.6d}$$

From equation (2.4), by virtue of equations (2.5) and (2.6), we can obtain

$$U_{ijkh} l^i = Q_{jkh}, U_{ijkh} g^{jj} = 0, \tag{2.7} a$$

$$U^*_{ij} = U_{ijkh} g^{kh} = {}^1A_{ij} + {}^4A_{ij}, U^*_{ij} + U^*_{ji} = 0, \tag{2.7} b$$

$$U_{ijkh} - U_{jhik} = \{{}^4A_0 (D_{(2)} - 2D_{(3)}) + {}^3A_0 D_{(3)} - D_{(1)} ({}^1A_0 + {}^2A_0)\} (m_k n_h - m_h n_k) \cdot (m_i n_j - m_j n_i). \tag{2.7} c$$

In general, $m_k n_h \neq m_h n_k$, therefore from equation (2.7) c we can obtain

Theorem 2.1

In a three-dimensional Finsler space F^3 , the necessary and sufficient condition for the tensor U_{ijkh} to be symmetric in k and h is given by ${}^4A_0 (D_{(2)} - 2D_{(3)}) + {}^3A_0 D_{(3)} - ({}^1A_0 + {}^2A_0) D_{(1)} = 0$.

In F^3 , it is known that Matsumoto [6]

$$h_{ik} h_{jh} - h_{ih} h_{jk} = (m_i n_j - m_j n_i) (m_k n_h - m_h n_k),$$

Therefore, equation (2.7) c, can also be expressed as

$$\zeta_{(k,h)} [U_{ijkh} - \{{}^4A_0 (D_{(2)} - 2D_{(3)}) + {}^3A_0 D_{(3)} - ({}^1A_0 + {}^2A_0) D_{(1)}\} h_{ik} h_{jh}] = 0 \tag{2.8}$$

It is also known that the tensor D'_{ijkh} can be expressed as [9]:

$$D'_{ijkh} = (2 D_{(1)}^2 - D_{(2)} D_{(3)} + D_{(3)}^2) (h_{ik} h_{jh} - h_{ih} h_{jk}), \tag{2.9} a$$

Such that

$$D'_{ik} = D'_{ijkh} g^{jh} = (2 D_{(1)}^2 - D_{(2)} D_{(3)} + D_{(3)}^2) h_{ik} \tag{2.9} b$$

Therefore, from equations (2.8) and (2.9) a, we can obtain

$$\zeta_{(k,h)} [U_{ijkh} - \{{}^4A_0 (D_{(2)} - 2D_{(3)}) + {}^3A_0 D_{(3)} - ({}^1A_0 + {}^2A_0) D_{(1)}\} \cdot D'_{ijkh} (2 D_{(1)}^2 - D_{(2)} D_{(3)} + D_{(3)}^2)^{-1}] = 0 \tag{2.10}$$

Hence:

Theorem 2.2

In a three-dimensional Finsler space F^3 , curvature tensors U_{ijkh} and D'_{ijkh} are related by equation (2.10).

Also, from equations (2.8) and (2.9), one can obtain after simplification

$$C_{\zeta(k,h)} \{U_{ijkh}\} = -D'_{ijkh/0} + 2D_{(1)}^2 A_0(h_{ik}h_{jh} - h_{ih}h_{jk}) \quad (2.11)$$

Hence:

Theorem 2.3

In a three-dimensional Finsler space F^3 , curvature tensor U_{ijkh} and $D'_{ijkh/0}$ are related by equation (2.11).

From equation (2.4), one can also imagine that the tensor U_{ijkh} can be expressed as the difference of product of two tensors in the following form:

$$U_{ijkh} = S_{ik}T_{jh} - B_{jk}T_{ih}, \quad (2.12)$$

where we have assumed

$$S_{ik} = a_i m_k + b_i n_k \quad (2.13a)$$

and

$$T_{jh} = d_j m_h + e_j n_h \quad (2.13b)$$

such that the vectors a_i , b_i , d_j and e_j satisfy following relations:

$${}^1A_{ij} = a_i d_j - a_j d_i, \quad {}^2A_{ij} = a_i e_j - a_j e_i, \quad (2.14a)$$

$${}^3A_{ij} = b_i d_j - b_j d_i, \quad {}^4A_{ij} = b_i e_j - b_j e_i \quad (2.14b)$$

Using equations (2.6) a, b, c, d in equations (2.13) a, b, we can obtain

$$a_i d_j - a_j d_i = {}^1A_i m_j + {}^4A_i n_j - {}^1A_j m_i - {}^4A_j n_i + A (m_j n_i - m_i n_j), \quad (2.15) a$$

$$a_i e_j - a_j e_i = {}^4A_i m_j - {}^2A_i n_j - {}^4A_j m_i + {}^2A_j n_i + B (m_j n_i - m_i n_j), \quad (2.15) b$$

$$b_i d_j - b_j d_i = {}^4A_i m_j - {}^2A_i n_j - {}^4A_j m_i + {}^2A_j n_i + C (m_j n_i - m_i n_j), \quad (2.15) c$$

$$b_i e_j - b_j e_i = {}^2A_j m_i - {}^3A_j n_i - {}^2A_i m_j + {}^3A_i n_j + E (m_j n_i - m_i n_j) \quad (2.15) d$$

where

$$A = D_{(3)} ({}^1A_0 - {}^2A_0) - (D_{(1)} - D_{(4)}) {}^4A_0, \quad B = D_{(3)} ({}^4A_0 - {}^3A_0) + (D_{(1)} - D_{(4)}) {}^2A_0,$$

$$C = D_{(4)} ({}^1A_0 - {}^2A_0) + (D_{(2)} - D_{(3)}) {}^4A_0, \quad E = D_{(4)} ({}^4A_0 - {}^3A_0) - (D_{(2)} - D_{(3)}) {}^2A_0.$$

Using $a_i m^i = a'$, $a_i n^i = a^*$, $b_i m^i = b'$, $b_i n^i = b^*$ etc., in (2.15) a, we can get

$$a' d_j - a_j d' = ({}^1A_i m^i) m_j + ({}^4A_i m^i) n_j - {}^1A_j - A n_j \quad (2.16) a$$

$$a^* d_j - a_j d^* = ({}^1A_i n^i) m_j + ({}^4A_i n^i) n_j - {}^4A_j + A m_j \quad (2.16) b$$

By solving equations (2.16) a and (2.16) b, we can obtain values of a_j and d_j in the following form:

$$a_j = (a' d^* - a^* d')^{-1} [m_j \{({}^1A_i m^i) a^* - ({}^1A_i n^i - A) a'\} + n_j \{({}^4A_i m^i - A) a^* - ({}^4A_i n^i) a'\} - ({}^1A_j a^* - {}^4A_j a')] \quad (2.17) a$$

$$d_j = (a' d^* - a^* d')^{-1} [m_j \{({}^1A_i m^i) d^* - ({}^1A_i n^i + A) d'\} + n_j \{({}^4A_i m^i - A) d^* - ({}^4A_i n^i) d'\} + ({}^4A_j d' - {}^1A_j d^*)] \quad (2.17) b$$

Similarly, from equations (2.15) b, c, d, we can obtain

$$a_j = (e^* a' - a^* e')^{-1} [m_j \{ ({}^4A_i m^i) a^* - ({}^4A_i n^i + B) a' \} + n_j \{ ({}^2A_i n^i) a' - ({}^2A_i m^i + B) a^* \} - ({}^4A_j a^* + {}^2A_j a')], \quad (2.17)c$$

$$e_j = (e^* a' - a^* e')^{-1} [m_j \{ ({}^4A_i m^i) e^* - ({}^4A_i n^i + B) e' \} + n_j \{ ({}^2A_i n^i) e' - ({}^2A_i m^i + B) e^* \} - ({}^2A_j e' + {}^4A_j e^*)], \quad (2.17)d$$

$$b_j = (b' d^* - b^* d')^{-1} [m_j \{ ({}^4A_i m^i) b^* - ({}^4A_i n^i + C) b' \} + n_j \{ ({}^2A_i n^i) b' - ({}^2A_i m^i + C) b^* \} - ({}^2A_j b' + {}^4A_j b^*)], \quad (2.17)e$$

$$d_j = (b' d^* - b^* d')^{-1} [m_j \{ ({}^4A_i m^i) d^* - ({}^4A_i n^i + C) d' \} + n_j \{ ({}^2A_i n^i) d' - ({}^2A_i m^i + C) d^* \} - ({}^2A_j d' + {}^4A_j d^*)], \quad (2.17)f$$

$$b_j = (b' e^* - b^* e')^{-1} [m_j \{ ({}^2A_i n^i - E) b' - ({}^2A_i m^i) b^* \} + n_j \{ ({}^3A_i m^i - E) b^* - ({}^3A_i n^i) b' \} + ({}^3A_j b' + {}^2A_j b^*)], \quad (2.17)g$$

$$e_j = (b' e^* - b^* e')^{-1} [m_j \{ ({}^2A_i n^i - E) e' - ({}^2A_i m^i) e^* \} + n_j \{ ({}^3A_i m^i - E) e^* - ({}^3A_i n^i) e' \} + ({}^2A_j e^* + {}^3A_j e')] \quad (2.17)h$$

Remark

In equations (2.17), we have obtained two values each for vectors a_i, b_i, d_i and e_i . By equating these values, we get

$$({}^1A_i n^i - {}^4A_i m^i + A)(e^* a' - a^* e') = ({}^2A_i m^i + {}^4A_i n^i + B)(a' d^* - a^* d') \quad (2.18)a$$

$$({}^2A_i n^i + {}^3A_i m^i - E)(b' d^* - b^* d') = ({}^2A_i m^i + {}^4A_i n^i + C)(b' e^* - b^* e') \quad (2.18)b$$

$$({}^4A_i m^i - {}^1A_i n^i - A)(b' d^* - b^* d') = ({}^2A_i m^i + {}^4A_i n^i + C)(a' d^* - a^* d') \quad (2.18)c$$

$$({}^2A_i m^i + {}^4A_i n^i + B)(b' e^* - b^* e') = ({}^2A_i n^i + {}^3A_i m^i - E)(a' e^* - a^* e') \quad (2.18)d$$

Hence:

Theorem 2.4

In a three-dimensional Finsler space F^3 , if the tensor U_{ijkh} is expressed as in (2.12), the coefficients ${}^1A_i, {}^2A_i, {}^3A_i, {}^4A_i$ and constants $a^*, b^*, d^*, e^*, a', b', d'$ and e' satisfy equations (2.18) a, b, c, d.

U-Ricci Tensors

Let us assume that we define two tensors with the help of equation (2.4) as follows:

$$U_{ijkh} g^{jk} = U_{ih}^{(1)} = {}^1A_{ij} m^j m_h + {}^2A_{ij} m^j n_h + {}^3A_{ij} n^j m_h + {}^4A_{ij} n^j n_h \quad (3.1)$$

and

$$U_{ijkh} g^{jh} = U_{ik}^{(2)} = {}^1A_{ij} m^j m_k + {}^2A_{ij} n^j m_k + {}^3A_{ij} m^j n_k + {}^4A_{ij} n^j n_k \quad (3.2)$$

The tensors $U_{jh}^{(1)}$ and $U_{jh}^{(2)}$ are tensors similar to Ricci-tensors for curvature tensor P_{ijkh} , defined and studied by Shimada [11].

Using equations (2.6) a, b, c, d and ${}^1A_{ij} m^j = {}^1A^*_i, {}^2A_{ij} m^j = {}^2A^*_i, {}^3A_{ij} m^j = {}^3A^*_i, {}^4A_{ij} m^j = {}^4A^*_i, {}^1A_{ij} n^j = {}^1A'_i, {}^2A_{ij} n^j = {}^2A'_i, {}^3A_{ij} n^j = {}^3A'_i$ and ${}^4A_{ij} n^j = {}^4A'_i$, in equations (3.1) and (3.2), we get

$$U_{ih}^{(1)} = ({}^1A^*_i + {}^3A'_i) m_h + ({}^2A^*_i + {}^4A'_i) n_h \quad (3.3)$$

and

$$U_{ik}^{(2)} = ({}^1A^*_i + {}^2A'_i) m_k + ({}^3A^*_i + {}^4A'_i) n_k \quad (3.4)$$

where

$${}^1A^*_i = {}^1A_i - m_i ({}^1A_j m^j) - n_i ({}^4A_j m^j) + \{ D_{(3)} ({}^1A_0 - {}^2A_0) - 2{}^4A_0 D_{(1)} \} n_i, \quad (3.5)a$$

$${}^1A'_i = {}^4A_i - m_i({}^1A_j n^j) - n_i({}^4A_j n^j) - \{D_{(3)}({}^1A_0 - {}^2A_0) - 2{}^4A_0 D_{(1)}\} m_i, \quad (3.5)b$$

$${}^2A^*_i = {}^4A_i - m_i({}^4A_j m^j) + n_i({}^2A_j m^j) + \{D_{(3)}({}^4A_0 - {}^3A_0) + 2{}^2A_0 D_{(1)}\} n_i, \quad (3.5)c$$

$${}^2A'_i = -{}^2A_i - m_i({}^4A_j n^j) + n_i({}^2A_j n^j) - \{D_{(3)}({}^4A_0 - {}^3A_0) + 2{}^2A_0 D_{(1)}\} m_i, \quad (3.5)d$$

$${}^3A^*_i = {}^4A_i - m_i({}^4A_j m^j) + n_i({}^2A_j m^j) + \{D_{(1)}({}^2A_0 - {}^1A_0) + {}^4A_0(D_{(2)} - D_{(3)})\} n_i, \quad (3.5)e$$

$${}^3A'_i = -{}^2A_i - m_i({}^4A_j n^j) + n_i({}^2A_j n^j) - \{D_{(1)}({}^2A_0 - {}^1A_0) + {}^4A_0(D_{(2)} - D_{(3)})\} m_i, \quad (3.5)f$$

$${}^4A^*_i = -{}^2A_i + m_i({}^2A_j m^j) - n_i({}^3A_j m^j) + \{D_{(1)}({}^3A_0 - {}^4A_0) + {}^2A_0(D_{(2)} - D_{(3)})\} n_i, \quad (3.5)g$$

$${}^4A'_i = {}^3A_i + m_i({}^2A_j n^j) - n_i({}^3A_j n^j) + \{D_{(1)}({}^3A_0 - {}^4A_0) + {}^2A_0(D_{(2)} - D_{(3)})\} m_i \quad (3.5)h$$

Using equations (3.5) a, b, c, d, e, f, g, h in (equations (3.3) and (3.4), we can establish

$$C_{(l,h)}[U_{ih}^{(1)} - m_h({}^1A_i - {}^2A_i) - n_h({}^3A_i + {}^4A_i) - m_i n_h \{D_{(1)}({}^3A_0 + {}^4A_0) + D_{(2)}{}^2A_0 - D_{(3)}{}^3A_0\}] = 0 \quad (3.6)$$

and

$$C_{(l,h)}[U_{ih}^{(2)} - m_h({}^1A_i - {}^2A_i) - n_h({}^3A_i + {}^4A_i) - m_i n_h \{D_{(1)}({}^3A_0 + {}^4A_0) + D_{(2)}{}^2A_0 - D_{(3)}{}^1A_0\}] = 0 \quad (3.7)$$

These tensors also satisfy

$$U_{i0}^{(1)} = 0, U_{i0}^{(2)} = 0, U_{0h}^{(1)} = ({}^1A_0 - {}^2A_0) m_h + ({}^4A_0 + {}^3A_0) n_h, \quad (3.8)$$

$$U_{0h}^{(2)} = ({}^1A_0 - {}^2A_0) m_h + ({}^3A_0 + {}^4A_0) n_h = U_{0h}^{(1)} \quad (3.9)$$

Hence:

Theorem 3.1

In a three-dimensional Finsler space F^3 , the Ricci-tensors based on curvature tensor U_{ijkh} and defined by equations (3.3) and (3.4) satisfy equations (3.6), (3.7), (3.8) and (3.9).

Tensor V_{ijkh}

We now define a tensor V_{ijkh} , based on third order tensor D_{ijk} as follows:

Definition 4.1

In a three-dimensional Finsler space F^3 , based on third order symmetric tensor D_{ijk} , in analogy to tensor T_{ijkh} , we here define the tensor V_{ijkh} as follows:

$$V_{ijkh} = L D_{ijk/h} + l_h D_{ijk} + l_k D_{ijh} + l_j D_{ikh} + l_i D_{jkh} \quad (4.1)$$

Substituting value of D_{ijk} from equation (1.10), in (4.1), after some simplification by virtue of equation (1.4), we get

$$V_{ijkh} = \alpha_h m_i m_j m_k + \sum_{(l,j,k)} (\beta_h m_i m_j n_k - \gamma_h m_i n_j n_k) + \delta_h n_i n_j n_k \quad (4.2)$$

where

$$\alpha_h = L D_{(1)/h} + l_h D_{(1)} - 3 v_h D_{(3)}, \quad (4.3)a$$

$$\beta_h = L D_{(3)/h} + l_h D_{(3)} + 3 v_h D_{(1)}, \quad (4.3)b$$

$$\gamma_h = - \{L D_{(1)/h} + l_h D_{(1)} + v_h(3 D_{(2)} - 2D)\}, \quad (4.3)c$$

$$\delta_h = L D_{(2)/h} + l_h D_{(2)} - 3 v_h D_{(1)} \tag{4.3} d$$

From equation (2.2), we can get $V_{ijkh} = V_{jikh}$, i.e., V_{ijkh} is symmetric in first two indices i and j . Also

$$V_{ijkh} - V_{ijhk} = m_i m_j (\alpha_h m_k - \alpha_k m_h + \beta_h n_k - \beta_k n_h) + n_i n_j (\gamma_h m_k - \gamma_k m_h) + (m_i n_j + m_j n_i) (\beta_h m_k - \beta_k m_h + \gamma_h n_k - \gamma_k n_h), \tag{4.4}$$

$$(V_{ijkh} - V_{ijhk}) g^{kh} = 0, \tag{4.5} a$$

and

$$(V_{ijkh} - V_{ijhk}) g^{ij} = D (v_k m_h - v_h m_k) + (\beta_h n_k - \beta_k n_h). \tag{4.5} b$$

Hence:

Theorem 4.1

In a three-dimensional Finsler space F^3 , tensor V_{ijkh} is symmetric in i and j and satisfies equations (4.4), (4.5) a and (4.5) b.

In case we assume that V_{ijkh} is also symmetric in k and h , equation (4.4) on simplification will give

$(\alpha_h + \gamma_h) m_k - (\alpha_k + \gamma_k) m_h + \beta_h n_k - \beta_k n_h = 0$, which by virtue of equations (4.3) a, b, c leads to

$$\beta_h m^h + D v_h n^h = 0 \tag{4.6} a$$

or alternatively

$$\beta_h m^h + D v_{233} = 0. \tag{4.6} b$$

Hence:

Theorem 4.2

In a three-dimensional Finsler space F^3 , if the tensor V_{ijkh} is symmetric in k and h , equation (4.6) b, is satisfied.

From equations (4.3) a, b, c, d we can obtain

$$\begin{aligned} \alpha_0 &= \alpha_h l^h = L D_{(1)/0} + D_{(1)}, \beta_0 = \beta_h l^h = L D_{(3)/0} + D_{(3)}, \\ \gamma_0 &= \gamma_h l^h = -L D_{(1)/0} - D_{(1)}, \delta_0 = \delta_h l^h = L D_{(2)/0} + D_{(2)} \\ \alpha_h m^h &= L D_{(1)/h} m^h - 3 D_{(3)} v_{232}, \beta_h m^h = L D_{(3)/h} m^h + 3 D_{(1)} v_{232} \\ \gamma_h m^h &= -L D_{(1)/h} m^h - (3 D_{(2)} - 2 D) v_{232}, \delta_h m^h = L D_{(2)/h} m^h - 3 D_{(1)} v_{232}, \\ \alpha_h n^h &= L D_{(1)/h} n^h - 3 D_{(3)} v_{233}, \beta_h n^h = L D_{(3)/h} n^h + 3 D_{(1)} v_{233}, \\ \gamma_h n^h &= -L D_{(1)/h} n^h - (3 D_{(2)} - 2 D) v_{233}, \\ \delta_h n^h &= L D_{(2)/h} n^h - 3 D_{(1)} v_{233}, \end{aligned} \tag{4.7}$$

which show

Theorem 4.3

In a three-dimensional Finsler space

- $\alpha_0 + \gamma_0 = 0$, ii) $\beta_0 + \delta_0 = L D_{/0} + D$,

- $(\alpha_h + \gamma_h) m^h + D v_{232} = 0$, iv) $(\alpha_h + \gamma_h) n^h + D v_{233} = 0$,
- $(\beta_h + \delta_h) m^h = L D_{/h} m^h$ and vi) $(\beta_h + \delta_h) n^h = L D_{/h} n^h$.

From equation (3.2), we can observe that

$$V_{ijkh} g^{ij} = (\alpha_h - \gamma_h) m_k + (\beta_h + \delta_h) n_k$$

and

$$V_{ijkh} g^{kh} = (\alpha_h m^h + \beta_h n^h) m_i m_j + (\beta_h m^h - \gamma_h n^h) (m_i n_j + m_j n_i) - (\gamma_h m^h - \delta_h n^h) n_i n_j,$$

which by virtue of equation (4.3) can be expressed as

$$V_{ijkh} g^{ij} = \{2(L D_{(1)/h} + l_h D_{(1)}) - v_h (5 D_{(3)} - D_{(2)})\} m_k + (L D_{/h} + l_h D) n_k \quad (4.8)$$

and

$$V_{ijkh} g^{kh} = \{L D_{(1)/h} m^h + L D_{(3)/h} n^h - 3(D_{(3)} v_{232} - D_{(1)} v_{233})\} m_i m_j + \{L D_{(1)/h} m^h + L D_{(2)/h} n^h - 3 D_{(1)} v_{233} + (3D_{(2)} - 2D)v_{232}\} n_i n_j + \{L D_{(3)/h} m^h + L D_{(1)/h} n^h + 3 D_{(1)} v_{232} + (3 D_{(2)} - 2D)v_{233}\} \cdot (m_i n_j + m_j n_i) \quad (4.9)$$

Hence:

Theorem 4.4

In a three-dimensional Finsler space F^3 , V-tensor V_{ijkh} satisfies equations (4.8) and (4.9).

From equation (4.2), we can also obtain

$$V_{ijkh} l^i = 0, V_{ijkh} l^j = 0, V_{ijkh} l^k = 0, \quad (4.10a)$$

$$V_{ijkh} l^h = m_i m_j (\alpha_0 m_k + \beta_0 n_k) + (m_i n_j + m_j n_i) (\beta_0 m_k + \gamma_0 n_k) + n_i n_j \gamma_0 m_k, \quad (4.10b)$$

$$V_{ijkh} g^{ij} l^h = \beta_0 n_k, V_{ijkh} g^{ij} l^h m^k = 0, V_{ijkh} g^{ij} n^k l^h = \beta_0. \quad (4.10c)$$

Hence:

Theorem 4.5

In a three-dimensional Finsler space F^3 , tensor V_{ijkh} satisfies equations (4.10) a, b, c.

Equation (4.2) in a three-dimensional Finsler space F^3 , can also be expressed as

$$V_{ijkh} = \sum_{(l,j,k)} \{A_{hk} h_{ij} + B_{hk} n_i n_j\} \quad (4.11)$$

where A_{hk} and B_{hk} are second order tensors defined by

$$A_{hk} = \{(1/3) \alpha_h m_k + \beta_h n_k\} \quad (4.12)$$

and

$$B_{hk} = \{[(1/3) \delta_h - \beta_h] n_k - [(1/3) \alpha_h + \gamma_h] m_k\} \quad (4.13)$$

From equations (4.12) and (4.13) we get

$$A_{hk} + B_{hk} = (1/3) \delta_h n_k - \gamma_h m_k \quad (4.14a)$$

$$A_{ok} = \{(1/3) \alpha_0 m_k + \beta_0 n_k\}, \quad (4.14b)$$

$$B_{0k} = [\{(1/3) \delta_0 - \beta_0\} n_k - \{(1/3) \alpha_0 + \gamma_0\} m_k], \tag{4.14c}$$

which imply

$$A_{h0} + B_{h0} = 0 \tag{4.15a}$$

$$A_{0k} + B_{0k} = (1/3) \delta_0 n_k - \gamma_0 m_k \tag{4.15b}$$

Substituting the values of δ_0 and γ_0 from equation (4.7) we get

$$A_{0k} + B_{0k} = (L D_{(1)/0} + D_{(1)}) m_k + (1/3)(L D_{(2)/0} + D_{(2)}) n_k \tag{4.16}$$

Hence:

Theorem 4.6

In a three-dimensional Finsler space F^3 , second order tensors A_{hk} and B_{hk} satisfy equation (4.14), and A_{0k} and B_{0k} satisfy equation (4.16).

In any three-dimensional Finsler space F^3 , the V-tensor is defined by equation (4.11), which motivates us to give the following definition similar to T3-like Finsler space [8].

Definition 4.2

A Finsler space $F^n(n > 3)$, shall be called a V3-like Finsler space, if for arbitrary second order tensors A_{hk} and B_{hk} , satisfying $A_{h0} = 0, B_{h0} = 0$, its V-tensor, V_{ijkh} is non-zero and is expressed by an equation of the form (4.11).

Second Curvature Tensor

Let P_{hijk} , be the second curvature tensor in F^3 Matsumoto [6], then the Ricci tensors corresponding to them are defined as

$$P_{hk}^{(1)} = P^i_{hjk} = C_{k/h} - C^j_{hk/j} + P^i_{kr} C^r_{jh} - P^r_{hk} C_r \tag{5.1}$$

and

$$P_{hk}^{(2)} = P^i_{hjk} = C_{k/h} - C^j_{hk/j} + C^r_{kh} C_{r/0} - P^i_{hr} C^r_{kj} \tag{5.2}$$

These tensors are non-symmetric and satisfy

$$P_{h0}^{(1)} = 0, P_{h0}^{(2)} = 0, P_{0k}^{(1)} = C_{k/0} = P_k, P_{0k}^{(2)} = C_{k/0} = P_k \tag{5.3}$$

If we assume that the tensor $A_{hk} = P_{hk}^{(1)}$, we can obtain by virtue of equation (5.1) and $*V_{ij} = V_{ijkh} g^{kh}$, where $*V_{ij}$ is a symmetric tensor

$$*V_{kh} = 4P_{hk}^{(1)} + B_{hk} + 2 B_{hi} n^i n_k \tag{5.4}$$

From equation (5.4) on simplification, we get

$$C_{\zeta(h,k)} \{ B_{ki} (\delta^i_h + 2 n^i n_h) - 4 (C_{k/h} + P^s_{kr} C^r_{sh}) \} = 0, \tag{4.5}$$

which leads to

$$B_{k0} = B_{0i} (\delta^i_k + 2 n^i n_k) + 4P_k \tag{5.6}$$

From equation (5.6), we can easily obtain

$$(B_{k0} - B_{0k}) m^k = 4 C_{/0}, (B_{k0} - 3 B_{0k}) n^k = 0. \quad (5.7)$$

Hence:

Theorem 5.1

In a D3-like, three-dimensional Finsler space F^3 , if the tensor A_{hk} is given by $P_{hk}^{(1)}$, the tensor B_{k0} satisfies equations (5.6) and (5.7).

D-Reducible Finsler Space F^3

In a D-reducible Finsler space F^3 , the tensor D_{ijk} satisfies Rastogi [9]

$$D_{ijk} = (1/4) \sum_{(i,j,k)} \{h_{ij} D_k\} \quad (6.1)$$

Equation (6.1) is similar to the one defined and studied by Matsumoto [4] for a C-reducible Finsler space. Equation (6.1) implies that $D_{(1)} = 0$, $D_{(2)} = (3/4) D$, $D_{(3)} = (1/4) D$, therefore from equation (2.3), we shall have

$$Q_{ijk} = (1/4) [D_{/0} (3 n_i n_j n_k + m_i m_j n_k + m_j m_k n_i + m_k m_i n_j) - D h_0 (3 m_i m_j m_k + m_i n_j n_k + m_j n_k n_i + m_k n_i n_j)] \quad (6.2)$$

From equations (6.1) and (6.2), we can have

$$Q_{ijk} = D^{-1} D_{/0} D_{ijk} - (1/4) D h_0 (3 m_i m_j m_k + m_i n_j n_k + m_j n_k n_i + m_k n_i n_j) \quad (6.3)$$

Following Izumi [2], we here give following definition:

Definition 6.1

A three-dimensional Finsler space F^3 , for a constant λ , shall be called, a Q^* -Finsler space, if the tensor $Q_{ijk} = \lambda D_{ijk}$.

From equation (6.3), we can observe that for a Q^* -Finsler space F^3 , $D^{-1} D_{/0} = \lambda$ and

$$D h_0 (3 m_i m_j m_k + m_i n_j n_k + m_j n_k n_i + m_k n_i n_j) = 0, \quad (6.4)$$

which on simplification implies $h_0 = 0$. Hence:

Theorem 6.1

In a D-reducible Q^* -Finsler space $F^3, \lambda = (\log D)_{/0}$ and $h_0 = 0$.

From equations (6.1), we can obtain

$$D_{(1)} = 0, D_{(2)} = (3/4) D \text{ and } D_{(3)} = (1/4) D, {}^1A_i = -(3/4) D h_i, {}^2A_i = (1/4) D h_i,$$

$${}^3A_i = (3/4) D_{/1}, {}^4A_i = (1/4) D_{/1}, {}^1A_0 = -(3/4) D h_0, {}^2A_0 = (1/4) D h_0,$$

$${}^3A_0 = (3/4) D_{/0}, {}^4A_0 = (1/4) D_{/0}.$$

Using these values, we can obtain from equations (2.6) a, b, c, d

$${}^1A_{ij} = C_{(i,j)} \{ (1/4) D_{/1} n_j - (3/4) D h_i m_j - (1/4) D^2 h_0 m_j n_i \}, \quad (6.5a)$$

$${}^2A_{ij} = C_{(i,j)} \{ (1/4) D_{/1} m_j - (1/4) D h_i n_j - (1/8) D D_{/0} m_j n_i \}, \quad (6.5b)$$

$${}^3A_{ij} = C_{(i,j)} \{ (1/4) D_{/1} m_j - (1/4) D h_i n_j + (1/8) D D_{/0} m_j n_i \}, \quad (6.5c)$$

$${}^4A_{ij} = \zeta_{(i,j)} \{ (3/4) D_{/i} n_j - (1/4) D h_i m_j - (1/8) D^2 h_0 m_j n_i \}. \tag{6.5d}$$

Substituting from equations (6.5) a, b, c, d in equation (2.4), on simplification we can obtain the value of U_{ijkh} , which easily gives

$$U_{ijkh} - U_{ijhk} = (1/4) D D_{/0} (m_j n_i - m_i n_j) (m_h n_k - m_k n_h) \tag{6.6}$$

From equation (6.6) we can obtain

Theorem 6.2

In a three-dimensional D-reducible Finsler space F^3 , the symmetry of U_{ijkh} in k and h will imply $D_{/0} = 0$.

From equation (2.9) a, b for a D-reducible Finsler space F^3 , we get

$$D'_{ijkh} = - (D^2/8) (h_{ik} h_{jh} - h_{ih} h_{jk}) \tag{6.7 a}$$

and

$$D'_{ik} = - (D^2/8) h_{ik} \tag{6.7b}$$

while from equation (2.10), we get

$$\zeta_{(k,h)} [U_{ijkh} - (1/4) D D_{/0} (h_{ik} h_{jh} - h_{ih} h_{jk})] = 0 \tag{6.8}$$

From equations (6.7) a and (6.8) we can obtain

$$\zeta_{(k,h)} [U_{ijkh} + 2 D^{-1} D_{/0} D'_{ijkh}] = 0 \tag{6.9}$$

Hence:

Theorem 6.3

In a three-dimensional D-reducible Finsler space F^3 , tensors U_{ijkh} and D'_{ijkh} are related by equation (6.9).

In case of a D-reducible Finsler space, from equations (3.6) and (3.7), on simplification we can obtain

$$\zeta_{(i,h)} [U_{ih}^{(1)} - D_{/i} n_h + D \{ h_i m_h + (3/16) (D_{/0} - D h_0) m_i n_h \}] = 0 \tag{6.10}$$

and

$$\zeta_{(i,h)} [U_{ih}^{(2)} - D_{/i} n_h + D \{ h_i m_h - (3/8) D h_0 m_i n_h \}] = 0 \tag{6.11}$$

Hence:

Theorem 6.4

In a three-dimensional D-reducible Finsler space F^3 , tensors $U_{ih}^{(1)}$ and $U_{ih}^{(2)}$ satisfy equations (6.10) and (6.11).

In a three-dimensional D-reducible Finsler space F^3 , from equations (4.3) a, b, c, d, we can obtain

$$\alpha_h = - (3/4) D v_h, \beta_h = (1/4) (L D_{/h} + D h_h), \gamma_h = - (1/4) D v_h, \delta_h = 3 \beta_h,$$

$$\alpha_0 = 0, \beta_0 = (1/4) (L D_{/h} 1^h + D), \gamma_0 = 0, \delta_0 = 3 \beta_0,$$

$$\alpha_h m^h = - (3/4) D v_{2j32}, \beta_h m^h = (1/4) (L D_{/h} m^h), \gamma_h m^h = (1/3) \alpha_h m^h,$$

$$\delta_h m^h = 3 \beta_h m^h, \alpha_h n^h = - (3/4) D v_{233}, \beta_h n^h = (1/4)(L D_{/h} n^h),$$

$$\gamma_h n^h = (1/3) \alpha_h n^h, \delta_h n^h = 3 \beta_h n^h.$$

From equations (4.8), (4.9) and (4.10) b, for a D-reducible Finsler space F^3 , we get

$$V_{ijkh} g^{ij} = (L D_{/h} + D I_h) n_k - (3/2) D m_k, \quad (6.12)a$$

$$V_{ijkh} g^{kh} = (1/4) [(L D_{/h} n^h - 3 D v_{232}) m_i m_j + (3 L D_{/h} n^h + D v_{232}) n_i n_j + (L D_{/h} m^h + D v_{233}) (m_i m_j + n_i n_j)], \quad (6.12)b$$

$$V_{ijkh} l^h = (1/4)(L D_{/0} + D) \sum_{(i,j,k)} \{ m_i m_j n_k \} \quad (6.12)c$$

Hence:

Theorem 6.5

In a three-dimensional D-reducible Finsler space F^3 , tensor V_{ijkh} satisfies equations (6.12) a, b, c.

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